

# Note on "Anomalous Hydrodynamic Drafting of Interacting Flapping Flags"

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We make remarks on Ristroph and Zhang's [*Phys. Rev. Lett.* 101, 194502 (2008)] paper.

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Ristroph and Zhang just showed that [1], inverted drafting could be produced by flexible flags (of which flapping in front reduces fluid forces). As reported in their experiments on *schooling* flapping flags, they found that it is the leader of a group who enjoys a significant drag reduction (of up to 50%), while the downstream flag suffers a drag increase. Thus, they remarked that this counterintuitive inverted drag relationship is rationalized by dissecting the mutual influence of shape and flow in determining drag [1].

The present author, however, likes to issue some comments here about Ristroph and Zhang's experimental procedure and data. Firstly, the technique for measurements of the drag on a flag originated from [2] (cf. Ref. 15: *The flagpole is fixed to a cantilever which bows slightly ( $< 0.5\text{mm}$ ) under the fluid forcing of the flag. The deflection is measured optically* in [1]; relating support deflection to drag). However, is the net deflection only due to the net force? We know that once there is a net moment (torque included) applied to one end of a (flat) cantilever there will be also a net deflection [3]. Meanwhile, can the net **thrust** (generated) [4] be distinguished from the net drag in [1-2]?

Note that in verifying the drag measuring technique [2] the flow is presumed to be described by the inviscid, incompressible Euler equations [2] by neglecting possible effects of flow compressibility due to thickness variations in the soap film (skin friction or *viscous drag* was also neglected [2]). The latter (compressibility [5] as well as variations of film's thickness [6]) is still being argued [7] and there is a possibility that **Marangoni** flows (due to surface-tension-driven mechanical instability) occur [8] and will induce errors to the experimental measurements reported in [1]. According to [5], the sound speed in the film is  $\sqrt{2E_m/\rho_f d_f} \approx 4.3\text{ m/s}$  ( $E_m$  : Marangoni elasticity modulus ( $= 0.088\text{ N/m}$  [5]),  $\rho_f$  : fluid density,  $d_f (= 4.7\mu\text{m}$  [1]) : film thickness). The flow velocity ( $U$ ) is  $2\text{ m/s}$  in [1] and then the Mach number is around  $0.46$  which means the flow is compressible. The theoretical validation for those relating bending deflection to drag in [1-2] which was based on the incompressible flow thus should be checked again and possibly be modified.

Meanwhile as noted in Ref. 15 of [1], the *minute* bending of the support is optically measured by the same procedure in [2]. How can *The flat cantilever suppresses lateral motion, and streamwise force fluctuations are damped by a viscous dashpot attached to the beam* (cf. Ref. 15 in [1]) be still valid during large-amplitude flapping motions? We know that the identification of an (equilibrium) neutral axis (n.a.) [3] is crucial to the judgement of the bending behavior of a beam (positive or negative deflection?). Can this n.a. be easily found for largely fluctuating support? The other argue is about the fixed width ( $9.5\text{ cm}$ ) of the planar water tunnel in [1]? How about the **interference** or induced **blockage** [9] (e.g. 4th. flag in Fig. 4a of [1]) between bodies, wakes, and edges of films when large-amplitude flapping occurs (cf. Fig. 1)? What happens once the width (soapy water descends in-between) increases or decreases a little? Is  $9.5\text{ cm}$  optimal for the conclusion made in [1] (especially for smaller separations :  $G/L$ , cf. [1] for  $G/L$  details, where reported significant anomalous inverted drafting appears; cf. Figs. 2 and 4 in [1]). Finally, considering the limitation for the optical resolution in [1] : can the total drag and drag increment be measured separately for  $G/L = 0$ ? How to calibrate the elastic response of the support instantaneously? These issues will dominate the conclusion made in [1]. *Acknowledgements.* The author thanks Ms. Chu Mary and Hsieh Jiu Xiang for their support.

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- [1] L. Ristroph and J. Zhang, *Phys. Rev. Lett.* **101**, 194502 (2008).
  - [2] S. Alben, M. Shelley, and J. Zhang, *Nature* **420**, 479 (2002); *Phys. Fluids* **16**, 1694 (2004).
  - [3] D.H. Hodges, *AIAA J.* **41**, 1131 (2003). A.E.H. Love, *A Treatise on the Mathematical Theory of Elasticity* (Dover, New York, 1944).
  - [4] J. Liao, *Phil. Trans. R. Soc. B* **362**, 1973 (2007).
  - [5] M. Beizaie and M. Gharib, *Exp. Fluids* **23**, 130 (1997).
  - [6] B.J. Carroll and J. Lucassen, *Chem. Eng. Sci.* **28**, 23 (1973).
  - [7] M.A. Rutgers, X.L. Wu, and W.B. Daniel, *Rev. Sci. Instrum.* **72**, 3025 (2001).
  - [8] V.A. Nierstrasz and G. Frens, *J. Colloid Interf. Sci.* **234**, 162 (2001).
  - [9] S.W.D. Wolf, *Prog. Aerospace Sci.* **31**, 85 (1995). U. Ganzer, *Prog. Aerospace Sci.* **22**, 81 (1985). K. Pettersson and A.

Rizzi, Prog. Aerospace Sci. **44**, 295 (2008).

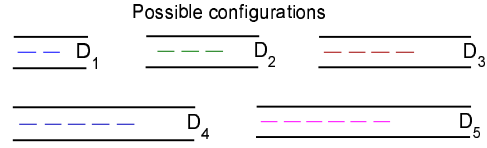


Fig. 1 Schematic set-up for different drag-force baseline calibration of tandem flags.  $D_0$  is for an isolated flag [1]. Ristroph and Zhang neglected the differences in [1] by presuming a universal  $D_0$  for the normalization and baseline comparison even the corresponding configurations are different (say, two and six tandem flapping flags).